

# Detecting Chameleons with Atom Interferometry

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## Outline:

Chameleon screening

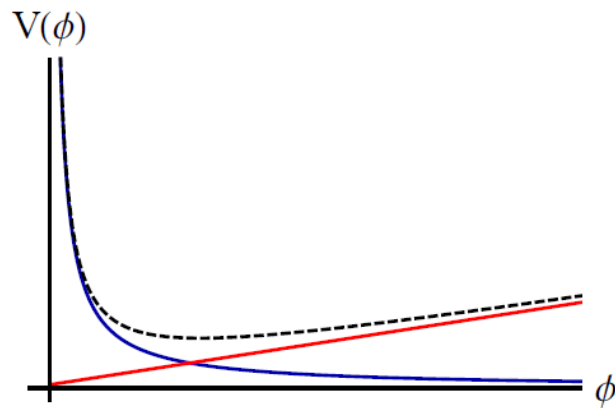
Atom interferometry

New constraints on chameleons

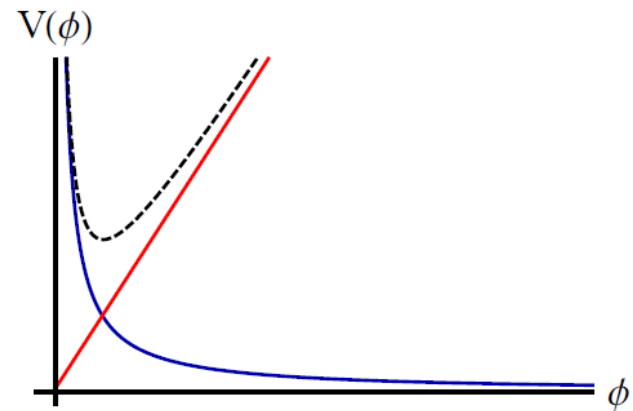
# The Chameleon - Varying Mass

The mass of the chameleon changes with the environment  
Field is governed by an effective potential

$$V_{\text{eff}} = \frac{\Lambda^5}{\phi} + \frac{\phi}{M} \rho$$



Low density



High density

**Warning:** Non-renormalisable theory, no protection from quantum corrections (But see [arXiv:1604.06051](https://arxiv.org/abs/1604.06051))

# The Scalar Potential

$$\phi = \phi_{\text{bg}} - \lambda_A \frac{1}{4\pi R_A} \frac{M_A}{M} \frac{R_A}{r} e^{-m_{\text{bg}} r}$$

$$\lambda_A = \begin{cases} 1, & \rho_A R_A^2 < 3M\phi_{\text{bg}} \\ 1 - \frac{S^3}{R_A^3} \approx 4\pi R_A \frac{M}{M_A} \phi_{\text{bg}}, & \rho_A R_A^2 > 3M\phi_{\text{bg}} \end{cases}$$

This determines how responsive an object is to the chameleon field

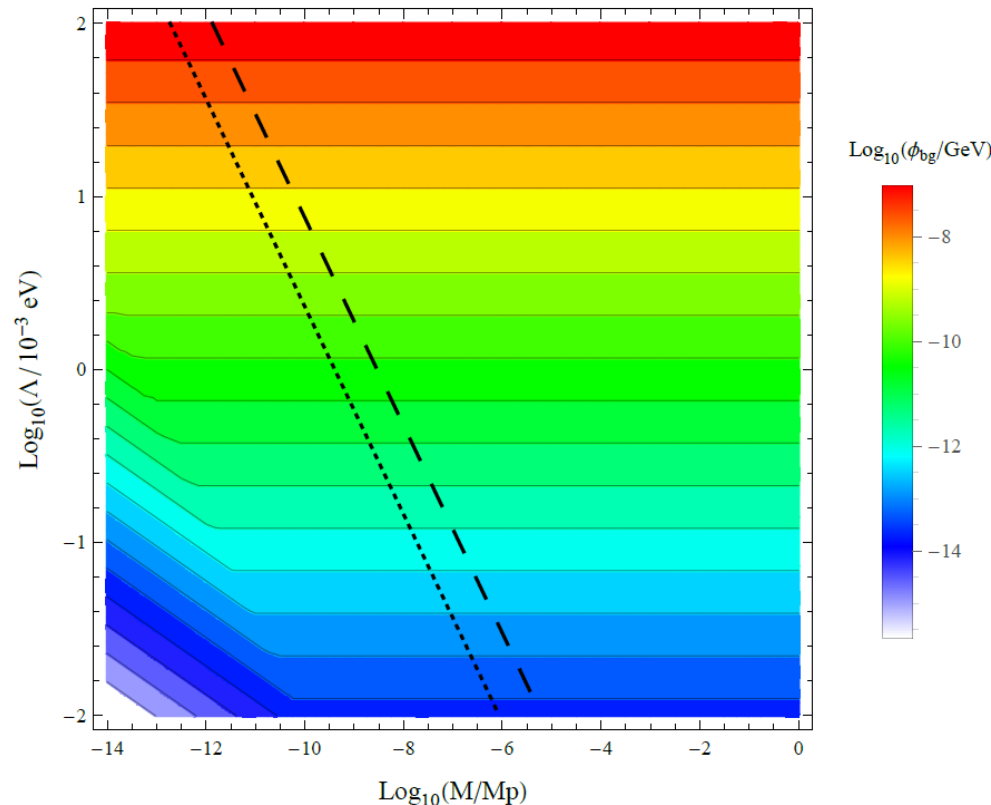
When  $m_{\text{bg}} r$  is small the ratio of the acceleration of a test particle due to the chameleon and gravity is:

$$\frac{a_\phi}{a_N} = \frac{\partial_r \phi}{M} \frac{r^2}{GM_A} = 3\lambda_A \left( \frac{M_P}{M} \right)^2$$

# Why Atom Interferometry?

In a spherical vacuum chamber, radius 10 cm, pressure  $10^{-10}$  Torr

Atoms are unscreened above black lines  
(dashed = caesium, dotted = lithium)

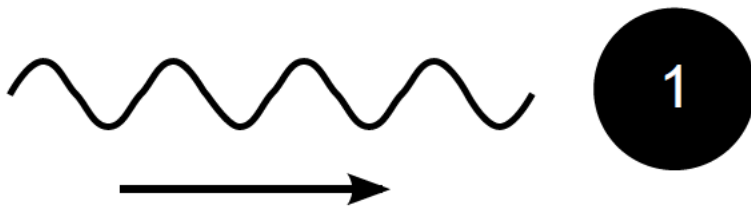


CB, Copeland, Hinds. (2015)

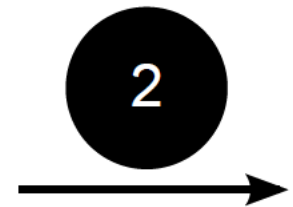
# What is Atom Interferometry?

An interferometer where the wave is made of atoms

Atoms can be moved around by absorption of laser photons

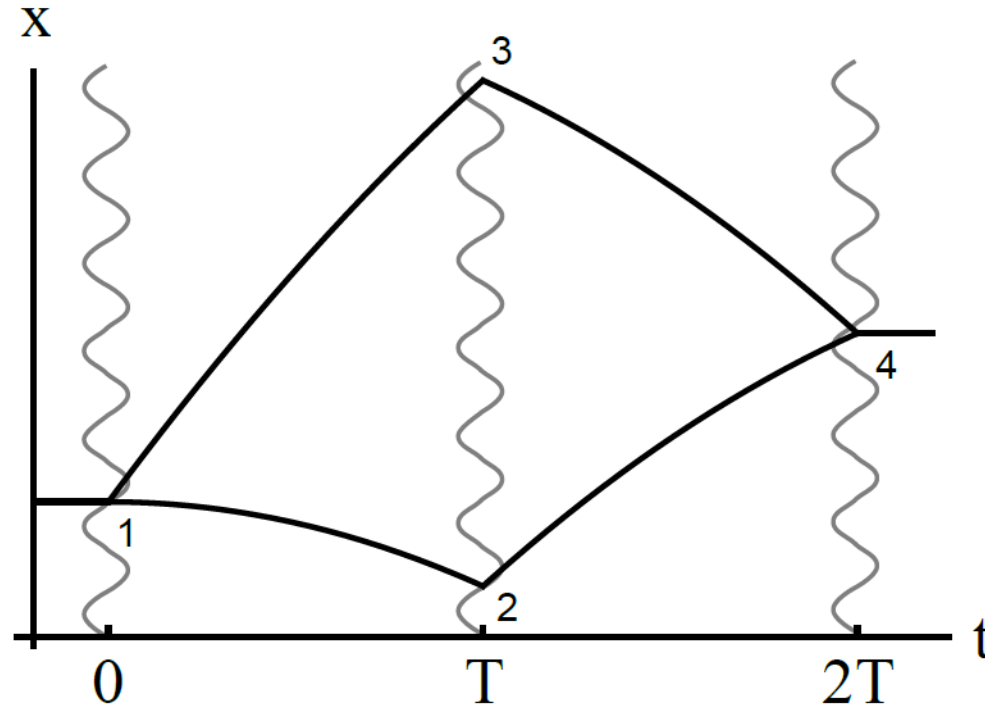


Photon Momentum =  $k$   
Atom in ground state



Atom in excited state  
with velocity =  $V$

# An Atom Interferometer



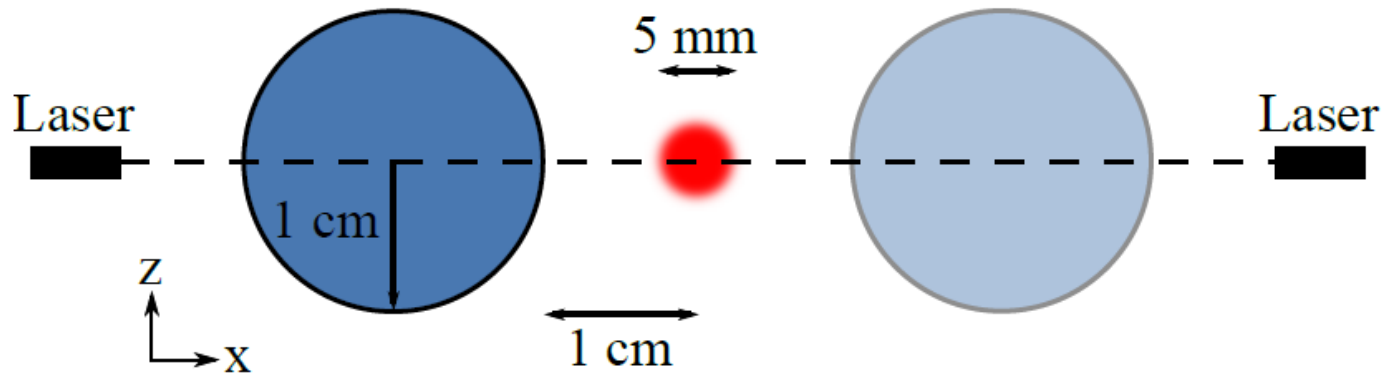
Probability measured in excited state at output

$$P = \cos^2 \left( \frac{kaT^2}{2} \right)$$

# Atom Interferometry for Chameleons

The walls of the vacuum chamber screen out any external chameleon forces

Macroscopic spherical mass (blue), produces chameleon potential felt by cloud of atoms (red)

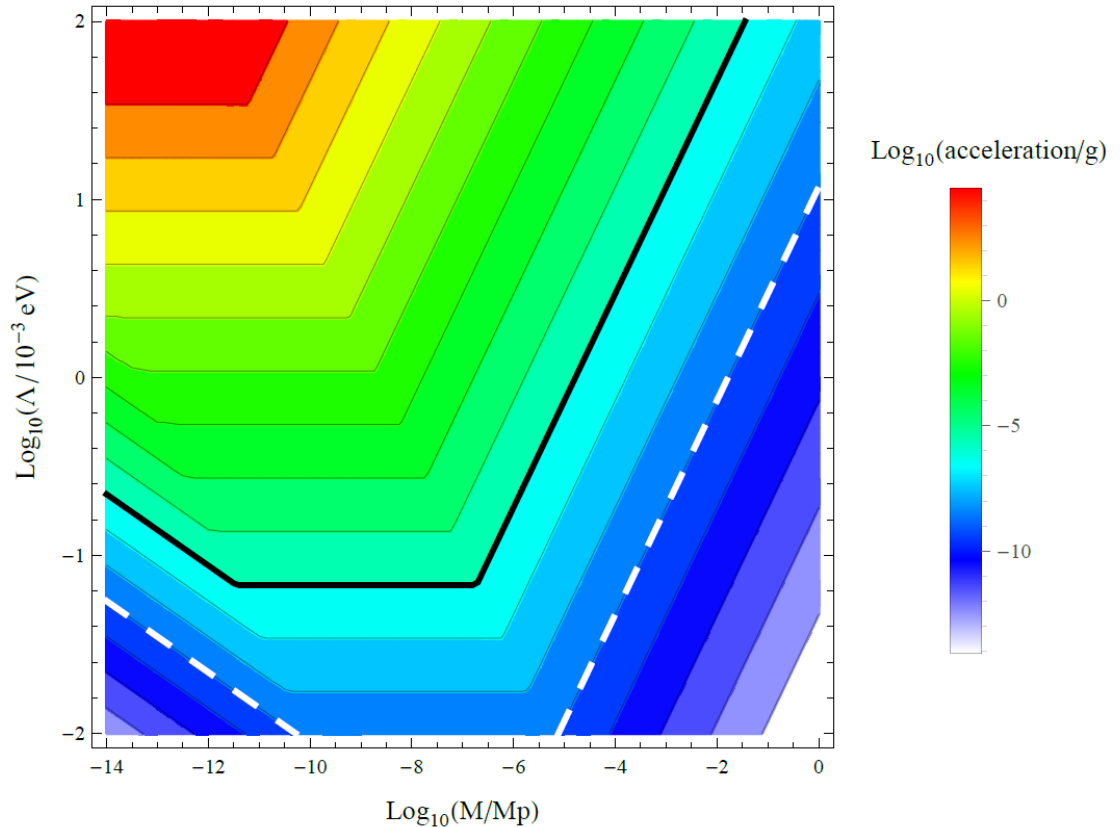


# Proposed Sensitivity

Systematics: Stark effect, Zeeman effect, phase shifts due to scattered light, movement of beams

All negligible at  $10^{-6}$  g sensitivity (solid black line)

Controllable down to  $10^{-9}$  g (dashed white line)



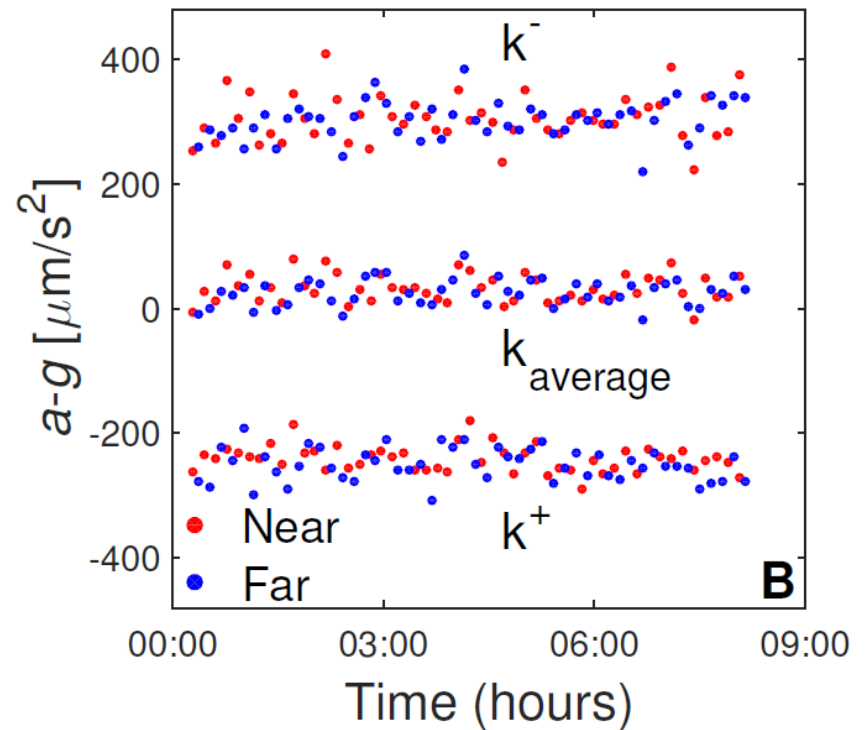
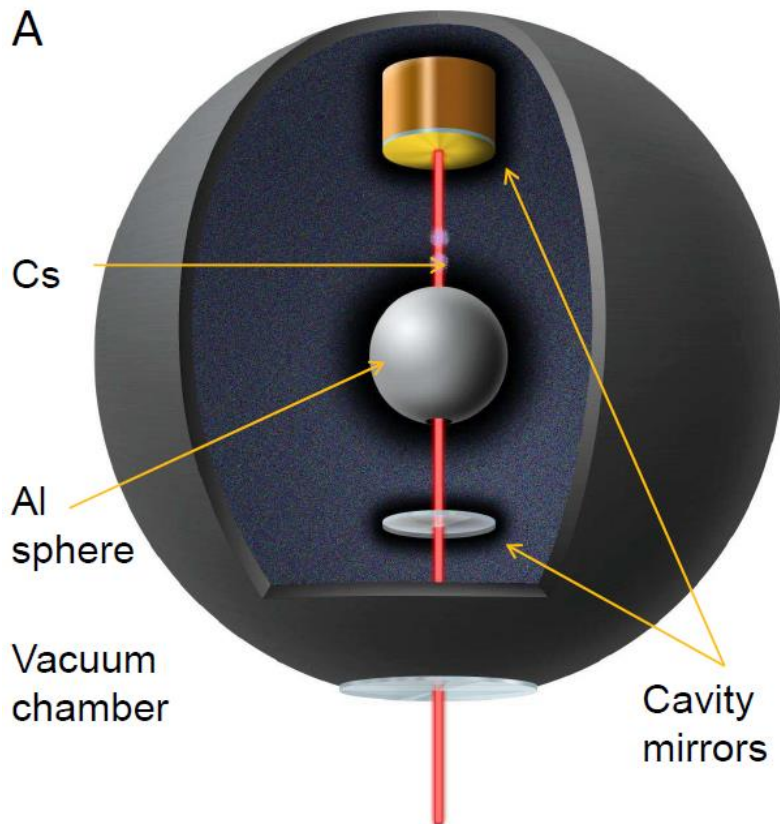
CB, Copeland, Hinds. (2015)

For numerical estimates see: Schlögel, Clesse, Füzfa (2015). Elder et al. (2016).

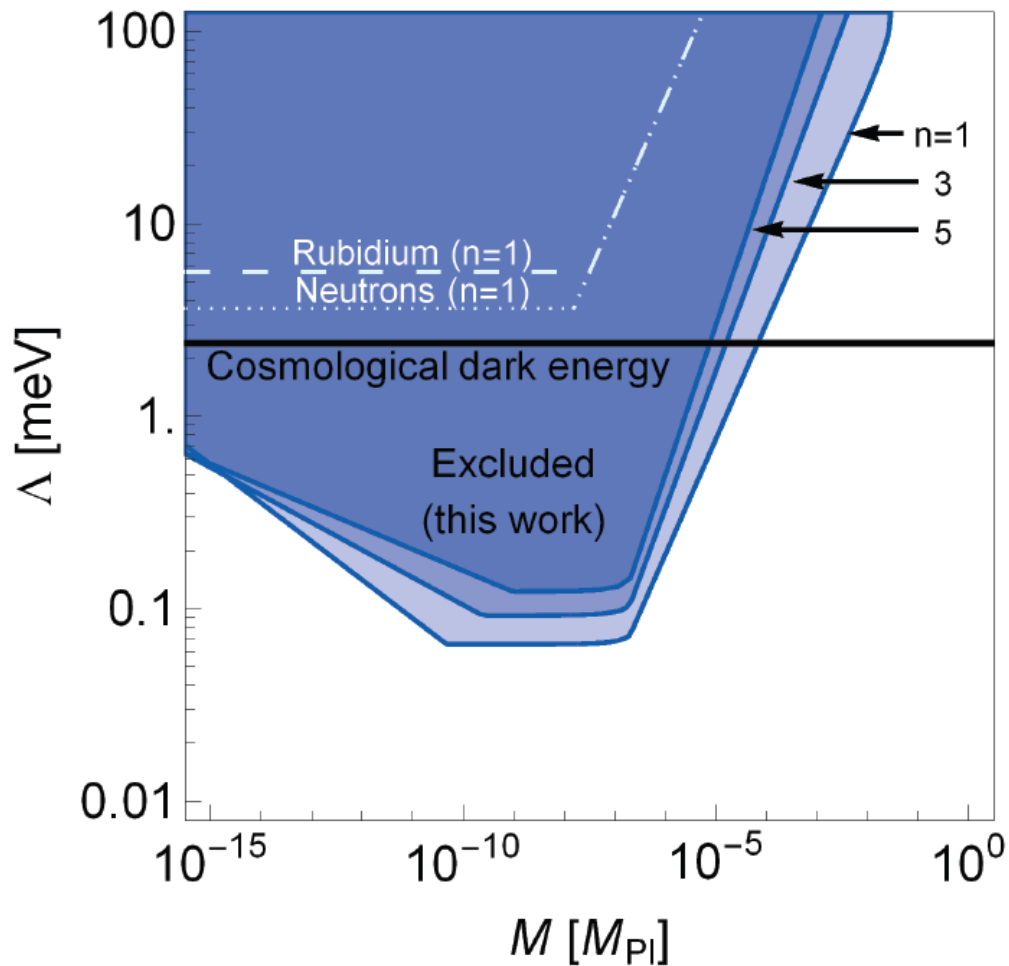


# Berkley Experiment

Using an existing set up with an optical cavity  
The cavity provides power enhancement, spatial filtering, and a precise beam geometry



# Berkley Experiment

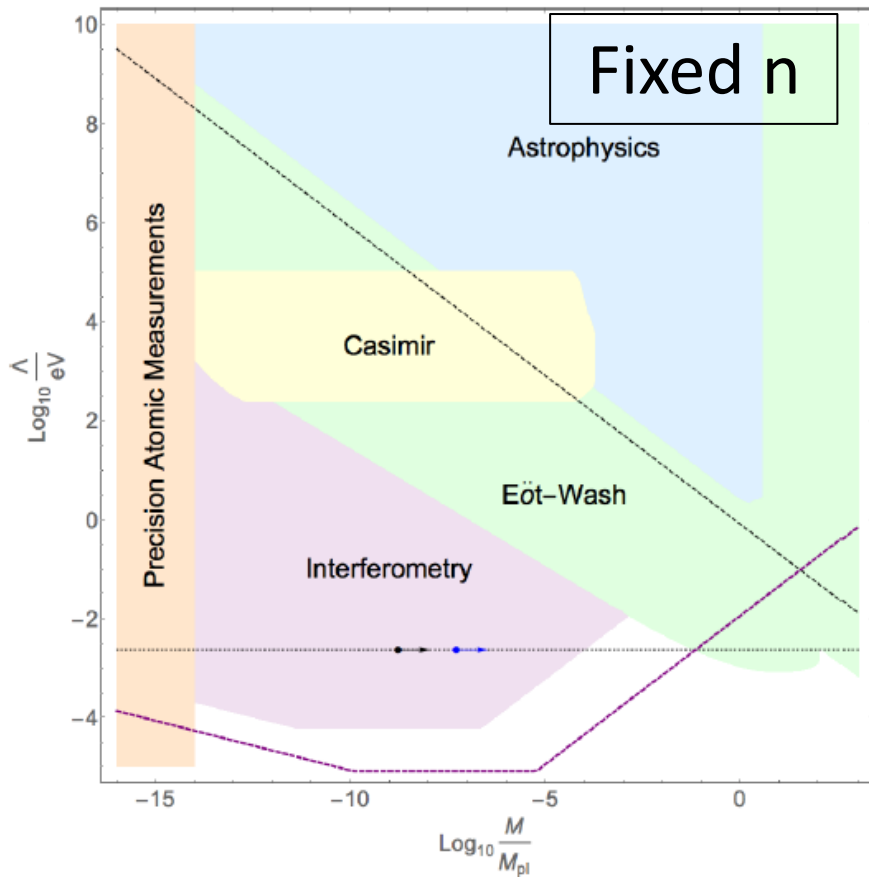


Hamilton et al. (2015)

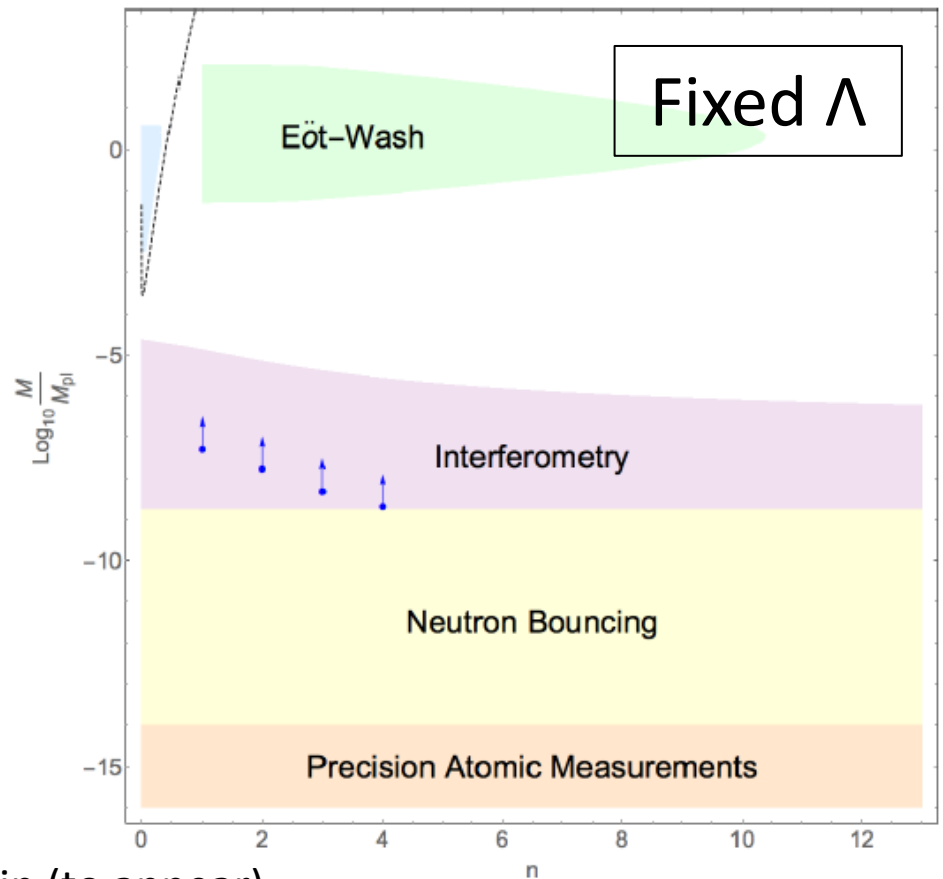
See also: Neutron interferometry experiments: Lemmel et al. 2015

Optically levitated microspheres: Rider et al. 2016

# Combined Constraints



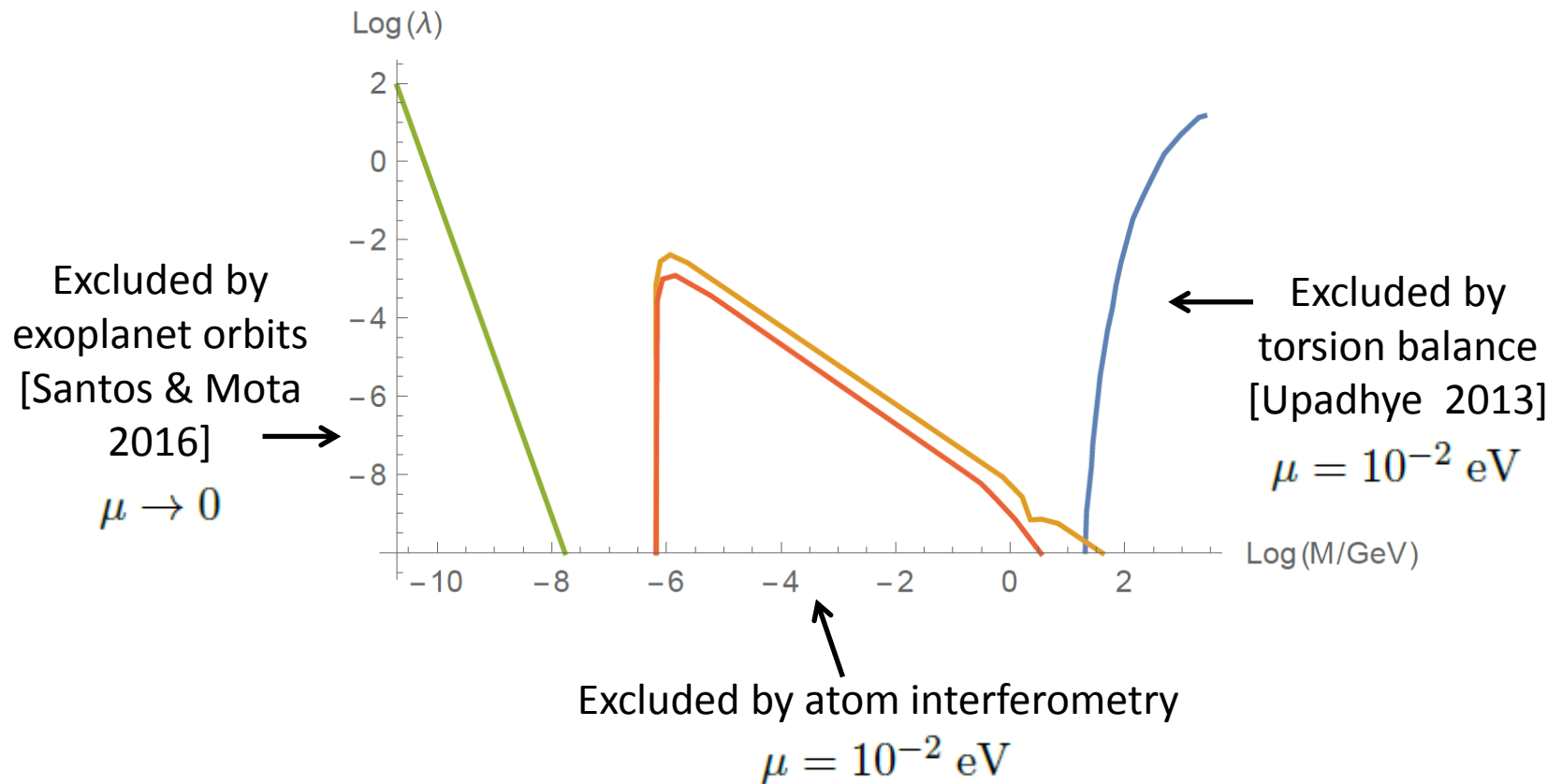
$$V(\phi) = \Lambda^4 + \frac{\Lambda^{4+n}}{\phi^n}$$



CB, Sakstein (to appear)

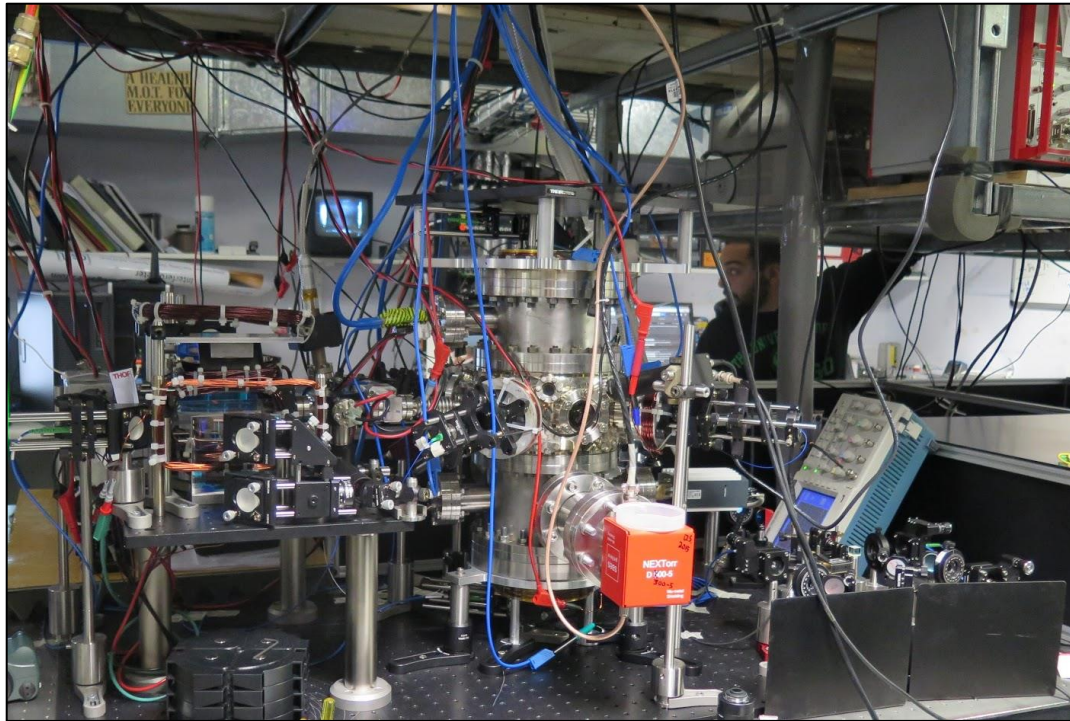
# Symmetron Constraints

$$V_{\text{eff}}(\phi) = \frac{1}{2} \left( \frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{1}{4} \lambda \phi^4$$



# Imperial Experiment

Development underway at the Centre for Cold Matter,  
Imperial College



Experiment rotated by 90 degrees from the Berkeley  
experiment, so that no sensitivity to Earth's gravity

# Summary

Screening mechanisms hide dark energy from standard searches for fifth forces

- Can still be detected in suitably designed experiments
- Atom interferometry a particularly powerful technique

Prospects for ruling out chameleon dark energy within a few years

Shape dependence of the screening mechanism could lead to important enhancements in sensitivity

