

Killing vectors

IF $\boxed{\mathcal{L}_{\vec{V}} \tilde{g} = 0}$, then $\vec{V}$ is called a Killing vector field.
 $\nearrow$ metric

Generates symmetries i.e. $\tilde{g}$ doesn't change along $\vec{V}$

Note: $(\mathcal{L}_{\vec{V}} g)_{\alpha\beta} = V^\gamma \nabla_\gamma g_{\alpha\beta} + g_{\alpha\beta} \nabla_\gamma V^\gamma + g_{\alpha\gamma} \nabla_\beta V^\gamma + g_{\gamma\beta} \nabla_\alpha V^\gamma$

$$= \boxed{\nabla_\alpha V_\beta + \nabla_\beta V_\alpha = 0} \quad \text{or } V_{(\alpha;\beta)} = 0$$

Killing's equation

Example: $ds^2 = d\theta^2 + \sin^2\theta d\phi^2$

$\frac{\partial}{\partial \phi}$ is a Killing vector i.e. metric independent of ϕ

Constants of the motion

Suppose $\vec{u}$ is tangent of geodesic. $\vec{V}$ is Killing Field

$$\nabla_{\vec{u}} \vec{u} = 0$$

Then $\nabla_{\vec{u}} (\vec{u} \cdot \vec{V}) = u^\alpha \nabla_\alpha (u^\beta V_\beta) = u^\alpha u^\beta \nabla_\alpha V_\beta$

$= 0$ because $\nabla_\alpha V_\beta$ is antisymmetric

so $\boxed{\vec{u} \cdot \vec{V} \text{ is constant along geodesic}}$

Killing vectors in flat spacetime

$$V_{\alpha,\beta} + V_{\beta,\alpha} = 0$$

↓

10 equations

⇒ 10 KVs

$$\Gamma = 0$$

$$g_{\mu\nu} = \eta_{\mu\nu}$$

(i) 4 constant KVs ⇒ 4 translations

$$V_{(\alpha)}^{\mu} = a_{(\alpha)}^{\mu} = \text{const}$$

← component of KV
↑ which KV

(ii) $V_{(\alpha)}^i = \epsilon^i_{km} x^m$, $V_{(\alpha)}^0 = 0$ ⇒ 3 Rotations

example $\vec{V}_{(x)} = x \frac{\partial}{\partial z} - z \frac{\partial}{\partial x}$

(iii) $V^{\mu(x)} = 2x^{\mu} x^k$ ⇒ 3 boosts

example $\vec{V}_{(x)} = -x \frac{\partial}{\partial t} - t \frac{\partial}{\partial x}$

(Why Boost? in coord transform language)

$$X'^{\alpha} = X^{\alpha} + \epsilon V^{\alpha}$$

$$t' = t - \epsilon x$$

$$x' = x - \epsilon t$$

$$y' = y$$

$$z' = z$$

Conserved Quantities

IF $\vec{V}$ is a KV, $J^\mu \equiv T^{\mu\nu} V_\nu$ is a conserved current

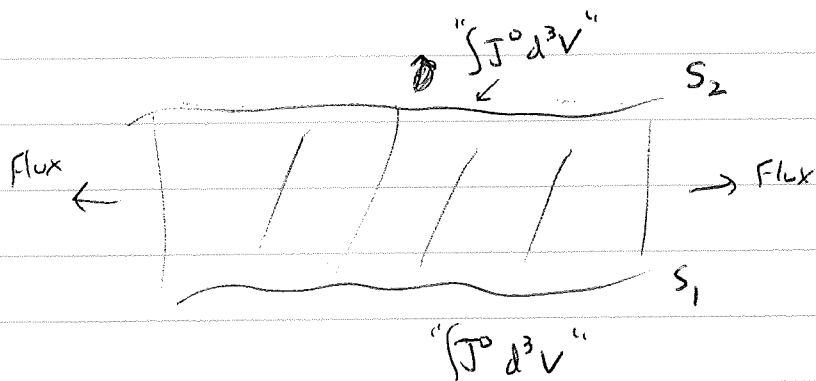
Why? $J^\mu_{;\mu} = (T^{\mu\nu} V_\nu)_{;\mu}$

$$= \underbrace{T^{\mu\nu}_{;\mu} V_\nu}_{D.T=0} + \underbrace{T^{\mu\nu}}_{\text{Sym}} \underbrace{V_{\nu;\mu}}_{\text{antisym}}$$

$$\Rightarrow \boxed{J^\mu_{;\mu} = 0}$$

Also get Conserved Charge

$$0 = \int J^\alpha_{;\alpha} \sqrt{g} d^4x = \overset{\text{sbkes/Gauss}}{\int J^\alpha d^3 \Sigma_\alpha}$$



IF J vanishes at spatial ∞
then

$$\int_{S_1} J^\alpha d^3 \Sigma_\alpha = \int_{S_2} J^\alpha d^3 \Sigma_\alpha$$

Conserved "currents/charges" in flat spacetime

(i) $\vec{V} = \frac{\partial}{\partial t}$ then $J^\alpha = T^{\alpha\beta} V_\beta = -T^{\alpha 0}$

$$Q = \int J^0 d^3x = \int \mathcal{E} d^3x \quad \text{energy}$$

$\vec{V} = \frac{\partial}{\partial x^i}$ $J^\alpha = -T^{\alpha i}$ $Q = \int J^0 d^3x = \int T^{0i} d^3x$
momentum

(ii) $\vec{V} = x \frac{\partial}{\partial z} - z \frac{\partial}{\partial x}$ $J^\alpha = x T^{\alpha y} - y T^{\alpha x}$

$$Q = \int (x T^{0y} - y T^{0x}) d^3x = "J_z"$$

angular momentum

(iii) boosts

$\vec{V} = x \frac{\partial}{\partial t} - t \frac{\partial}{\partial x}$ $J^\alpha = x T^{\alpha 0} - t T^{\alpha x}$

$$Q = \int (x T^{00} - t T^{0x}) d^3x$$

What is this?

define $X_{cm} = \frac{\int x T^{00}}{\int T^{00}}$ $V_{cm}^x = \frac{\int T^{0x}}{\int T^{00}}$

$$= \frac{P^x}{E}$$

So $Q = X_{cm} E - t P^x$

$$= E (X_{cm} - V_{cm}^x t)$$

$$= \text{const}$$

$\Rightarrow$ Uniform motion of CM of system

Linearized Gravity

- How to handle weak grav. Fields.

Assume $g_{\alpha\beta} = \eta_{\alpha\beta} + h_{\alpha\beta}$ $|h_{\alpha\beta}| \ll 1$

- Drop all h^2 terms (for solar system, $h \sim \frac{M}{R} \sim 10^{-6}$)

- Think of h as tensor field in flat space

- Raise / lower indices with η (but note $g^{\alpha\beta} = -h^{\alpha\beta} + \eta^{\alpha\beta}$)

Then: $\Gamma^\mu_{\alpha\beta} = \frac{1}{2} \eta^{\mu\nu} (h_{\alpha\nu,\beta} + h_{\beta\nu,\alpha} - h_{\alpha\beta,\nu})$ coord basis
 $= \frac{1}{2} h_{\alpha,\nu}^\mu + h_{\nu,\alpha}^\mu - h_{\alpha\beta}^{\mu\nu}$ (dropped 2nd order terms)

And: $R_{\mu\nu} = \Gamma^\alpha_{\mu\nu,\alpha} - \Gamma^\alpha_{\mu\alpha,\nu}$ [drop $\Gamma.\Gamma$ terms]

$$= \frac{1}{2} (h_{\mu,\alpha}^\alpha + h_{\nu,\alpha}^\alpha - h_{\mu\nu,\alpha}^\alpha - h_{,\alpha\nu}^\alpha)$$

where $h \equiv h^\alpha_\alpha$
 $= \eta^{\alpha\beta} h_{\alpha\beta}$

Equations simpler if we define $\bar{h}_{\mu\nu} \equiv h_{\mu\nu} - \frac{1}{2} \eta_{\mu\nu} h$

Then $G_{\mu\nu} = 8\pi T_{\mu\nu} \Rightarrow$
$$\boxed{-\bar{h}_{\mu\nu,\alpha}^\alpha - \eta_{\mu\nu} \bar{h}^{\alpha\beta}_{,\alpha\beta} + \bar{h}_{\mu,\alpha}^\alpha{}_{,\nu} + \bar{h}_{\nu,\alpha}^\alpha{}_{,\mu}} = 16\pi T_{\mu\nu}$$

Linearized Gravity

Classification

$$g_{\alpha\beta} = \eta_{\alpha\beta} + h_{\alpha\beta}$$

$$\text{but } g^{\alpha\beta} \neq \eta^{\alpha\beta} + h^{\alpha\beta}$$

$$\text{instead, } g^{\alpha\beta} = \eta^{\alpha\beta} - h^{\alpha\beta}$$

$$\text{so that } g^{\alpha\beta} g_{\beta\gamma} = \delta^\alpha_\gamma$$

In other words,

$$h^{\alpha\beta} = \eta^{\alpha\mu} \eta^{\beta\nu} h_{\mu\nu}$$

use η to raise/lower
 h

$$g^{\alpha\beta} \neq \eta^{\alpha\mu} \eta^{\beta\nu} g_{\mu\nu}$$

use g to raise/lower
 g

2 types of coord X Formation :

① Global:
Lorentz

$$X^{\mu'} = \Lambda^{\mu'}_{\mu} X^{\mu}$$

Then

$$g_{\bar{\alpha}\bar{\beta}} = (\eta_{\alpha\beta} + h_{\alpha\beta}) \Lambda^{\alpha}_{\bar{\alpha}} \Lambda^{\beta}_{\bar{\beta}}$$

$$= \eta_{\bar{\alpha}\bar{\beta}} + h_{\bar{\alpha}\bar{\beta}}$$

so $h_{\bar{\alpha}\bar{\beta}} = \Lambda^{\alpha}_{\bar{\alpha}} \Lambda^{\beta}_{\bar{\beta}} h_{\alpha\beta}$ like tensor in Flat spacetime

② Infinitesimal

$$X^{\mu}_{\text{new}}(P) = X^{\mu}_{\text{old}}(P) + \epsilon V^{\mu}(P) \quad \text{where } V^{\mu} \text{ is a vector field}$$

We know that $g_{\mu\nu}^{\text{new}} = g_{\mu\nu}^{\text{old}} + \epsilon \mathcal{L}_{\vec{V}} g_{\mu\nu}^{\text{old}}$

$$= g_{\mu\nu}^{\text{old}} - \epsilon [V^{\alpha} g_{\mu\nu,\alpha} + g_{\alpha\nu} V^{\alpha}_{,\mu} + g_{\alpha\mu} V^{\alpha}_{,\nu}]$$

$$\text{so } h_{\mu\nu}^{\text{new}} = h_{\mu\nu}^{\text{old}} - V_{\mu,\nu} - V_{\nu,\mu} + \mathcal{O}(h \cdot \epsilon)$$

Gauge X Formation

like $A_{\mu} \rightarrow A_{\mu} + \nabla_{\mu} \chi$ in $\mathbb{E} \times M$

$$\text{Now } h^{\text{new}} = h^{\text{old}} - 2V^{\alpha}_{,\alpha}$$

$$\begin{aligned} \text{so } \bar{h}^{\text{new}}_{\alpha\beta} &= h^{\text{new}}_{\alpha\beta} - \frac{1}{2} \eta_{\alpha\beta} h^{\text{new}} \\ &= h^{\text{old}}_{\alpha\beta} - V_{\alpha,\beta} - V_{\beta,\alpha} - \frac{1}{2} \eta_{\alpha\beta} (h^{\text{old}} - 2V^{\gamma}_{,\gamma}) \\ &= \bar{h}^{\text{old}}_{\alpha\beta} - V_{\alpha,\beta} - V_{\beta,\alpha} + \eta_{\alpha\beta} V^{\gamma}_{,\gamma} \end{aligned}$$

$$\text{Note that } \bar{h}^{\text{new},\beta}_{\alpha\epsilon} = \bar{h}^{\text{old},\beta}_{\alpha\epsilon} - V_{\alpha,\epsilon}{}^{\beta} - V_{\epsilon,\alpha}{}^{\beta} + \eta_{\alpha\epsilon} V^{\gamma}_{,\gamma}{}^{\beta}$$

$\xrightarrow{\text{cancel}} \quad \uparrow \quad \uparrow$
 $V^{\gamma}_{,\gamma\alpha}$

So can choose gauge transformation so that

$$\boxed{\bar{h}^{\text{new},\beta}_{\alpha\epsilon} = 0}$$

Lorentz
Gauge

by choosing $\vec{V}$

$$\text{s.t. } V_{\alpha\epsilon}{}^{\beta} = \bar{h}^{\text{old},\beta}_{\alpha\epsilon}$$

wave equation

Unique up to $V_{\alpha,\beta}{}^{\beta} = 0$ homogeneous solution.

So in Lorentz Gauge

$$\boxed{\begin{aligned} \bar{h}_{\mu\nu,\alpha}{}^{\alpha} &= -16\pi T_{\mu\nu} \\ \bar{h}_{\mu\nu}{}^{,\nu} &= 0 \end{aligned}}$$

Field
eqs.
in Lorentz
Gauge

$$\left(\begin{array}{l} \text{like } \square A^{\mu} = 4\pi J^{\mu} \\ A^{\mu}_{,\alpha} = 0 \end{array} \right)$$

In curvilinear coords, $\eta_{\mu\nu} \rightarrow g_{\mu\nu}$ ^{Flat}

$h_{\mu\nu}$ is tensor with respect to

so

$$\bar{h}_{\mu\nu;\alpha}{}^{\alpha} = -16\pi T_{\mu\nu}$$

$$\bar{h}_{\mu\nu}{}^{;\nu} = 0$$

New phenomenon

$$\text{In vacuum, } \bar{h}_{\mu\nu;\alpha}{}^\alpha = 0 \quad \bar{h}_{\mu\nu}{}^{;\nu} = 0$$

admits wave solutions

$$\text{example } \bar{h}_{xx} = \bar{h}_{xx}(t-z), \quad \bar{h}_{xy} = \bar{h}_{xy}(t-z), \quad \bar{h}_{yy} = \bar{h}_{yy}(t-z)$$

$$\bar{h}_{\mu 0} = \bar{h}_{\mu z} = 0$$

Gravitational waves

not present in Newtonian gravity.

$$\text{In general, } \bar{h}_{\mu\nu}(t, \underline{x}) = 4 \cdot \int \frac{T_{\mu\nu}(t - |\underline{x} - \underline{x}'|, \underline{x}') d^3 x'}{|\underline{x} - \underline{x}'|}$$

solution is

retarded time solution