

Geodesic Deviation

Equivalence principle: no such thing as uniform grav. Field
→ simply an accelerated frame.

Freely falling particle $\nabla_{\vec{u}} \vec{u} = 0$

Particle sitting on earth $\nabla_{\vec{u}} \vec{u} = \vec{a}$

$a^i = g$
upward.
in rest frame

So is there a gravitational field? Yes! tidal field, tidal forces

∴ Nearby geodesics initially parallel don't stay parallel.

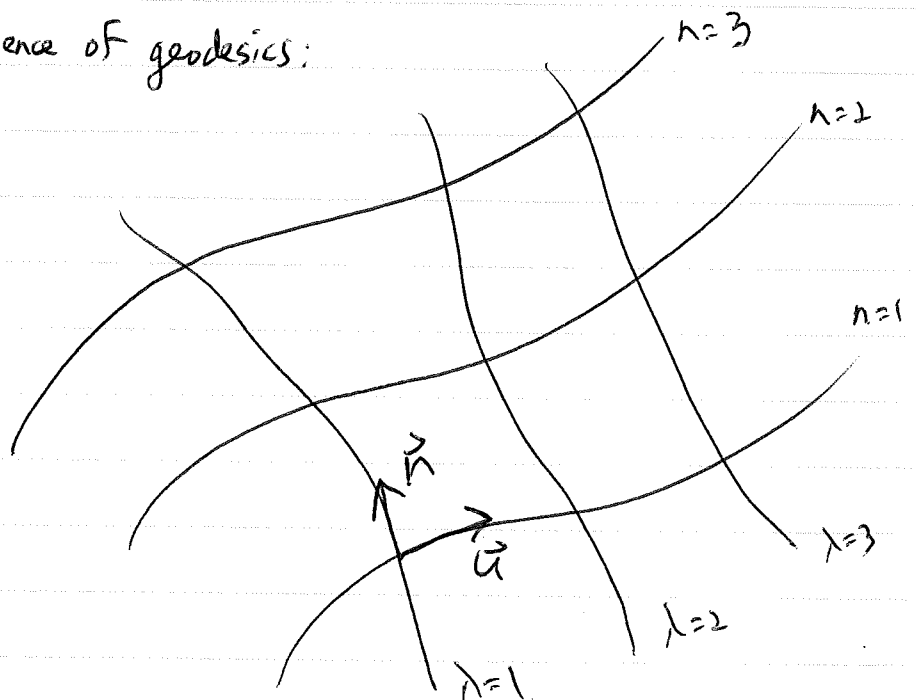
Derive geodesic deviation

Assume congruence of geodesics:

$$x^\alpha = x^\alpha(\lambda, n)$$

n = which geodesic

λ = affine param
on each geodesic



$\vec{n} = \frac{d}{d\lambda}$ connects points with same λ on different geodesics.

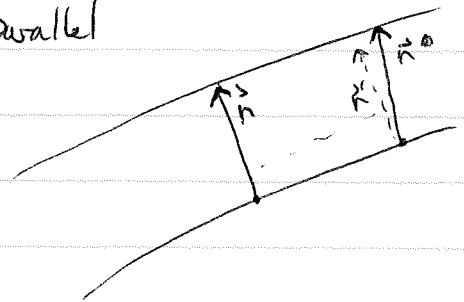
$\vec{u} = \frac{d}{d\lambda}$ tangent to geodesic.

By construction, $\vec{u}, \vec{n}$ form coord basis

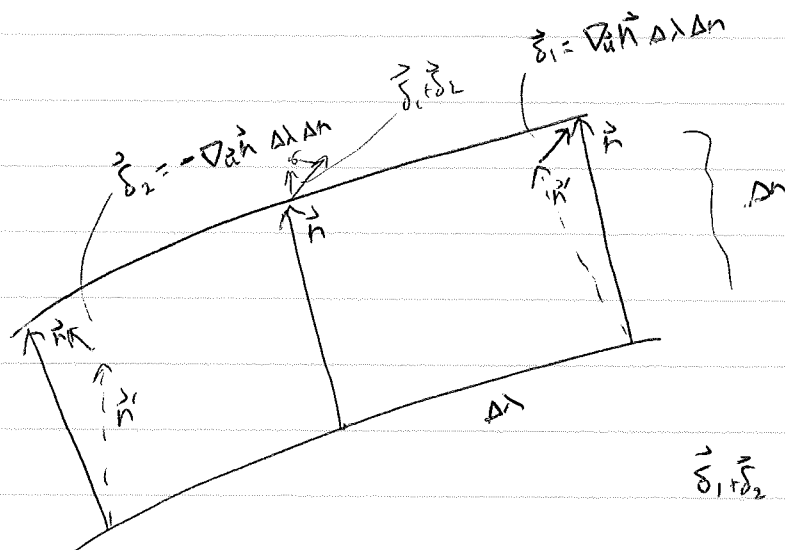
$$[\vec{u}, \vec{n}] = 0 = \nabla_{\vec{u}} \vec{n} - \nabla_{\vec{n}} \vec{u}$$

$\nabla_{\vec{u}} \vec{n}$ = how $\vec{n}$ changes along $\vec{u}$
= "velocity of nearby geodesic relative to this geodesic"

if $\nabla_{\vec{u}} \vec{n} = 0$, geodesics are parallel



$\nabla_{\vec{u}} \nabla_{\vec{u}} \vec{n}$ = "acceleration of nearby geodesic relative to this geodesic"



$$\vec{\delta}_1, \vec{\delta}_2 = \Delta \lambda^2 \nabla_{\vec{u}} \nabla_{\vec{u}} \vec{n}$$

So deviation from parallel geodesics $\vec{\delta}_1 + \vec{\delta}_2 = \nabla_{\vec{u}} \nabla_{\vec{u}} \vec{n} \Delta\lambda^2 \Delta n$

$$\text{or } \nabla_{\vec{u}} \nabla_{\vec{u}} \vec{n} = \frac{\vec{\delta}_1 + \vec{\delta}_2}{\Delta\lambda^2 \Delta n} = \text{relative acceleration}$$

$$\text{Now } \nabla_{\vec{u}} \nabla_{\vec{n}} \vec{n} = \nabla_{\vec{u}} \cancel{\nabla_{\vec{n}} \vec{u}} \quad \text{because } \nabla_{\vec{u}} \vec{n} = \nabla_{\vec{n}} \vec{u} \\ \text{ie. } [\vec{u}, \vec{n}] = 0$$

$$+ \tilde{R}(-, \vec{u}, \vec{u}, \vec{n})$$

$$- (\cancel{\nabla_{\vec{u}} \nabla_{\vec{n}} \vec{u}} - \nabla_{\vec{n}} \nabla_{\vec{u}} \vec{u} - \nabla_{[\vec{u}, \vec{n}]} \vec{u})$$

$$= \underbrace{\nabla_{\vec{n}} \nabla_{\vec{u}} \vec{u}}_0 \text{ geodesic} + \underbrace{\nabla_{[\vec{u}, \vec{n}]} \vec{u}}_{[\vec{u}, \vec{n}] = 0} + \tilde{R}(-, \vec{u}, \vec{u}, \vec{n})$$

$$\Rightarrow \boxed{\nabla_{\vec{u}} \nabla_{\vec{u}} \vec{n} = \tilde{R}(-, \vec{u}, \vec{u}, \vec{n})}$$

$$\text{or if we write } \nabla_{\vec{u}} \nabla_{\vec{u}} \vec{n} = \frac{D^2 \vec{n}}{d\tau^2}$$

$$\boxed{\frac{D^2 \vec{n}}{d\tau^2} = R^\alpha{}_{\beta\gamma\delta} u^\beta u^\gamma n^\delta}$$

Geodesic deviation

Curvature causes ^{nearby} geodesics to converge/diverge. Tidal Force

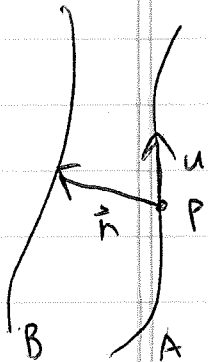
To keep two ^{slightly} separated particles both ^{parallel to each other}, must apply force.

More on tidal force

Assume observer ^A w/ 4-velocity $\vec{u}$.

Assume $\vec{n}$ is purely spatial as seen by A.

$$\vec{n} \cdot \vec{u} = 0$$



Then in rest frame of $\vec{u}$, $\frac{D^2 \vec{n}}{d\tau^2}$ = acceleration of Particle B measured by A.
 $= \frac{d^2 \vec{n}}{dt^2} = \Delta \vec{a}$

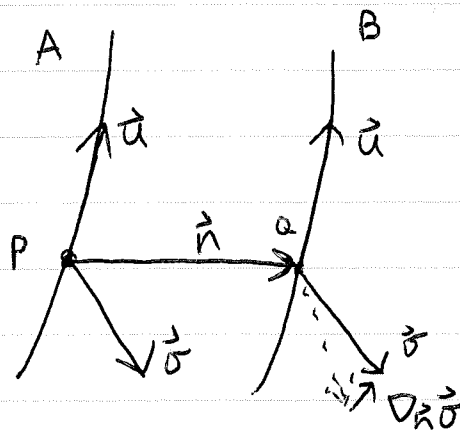
$$\begin{aligned} \frac{d^2 n^i}{dt^2} &= R^i{}_{\alpha\gamma\delta} u^\alpha u^\gamma n^\delta \\ &= R^i{}_{00j} n^j \quad \text{in rest frame} \end{aligned}$$

Sometimes we define $R^i{}_{0j0} = \mathcal{E}^i{}_j$

"Electric part of
Riemann tensor"
"Tidal Field"

$$\text{so } \Delta a^i = -\mathcal{E}^i{}_j n^j$$

Another effect of curvature:
Local Frame Dragging



Observers A & B separated by vector $\vec{n}$.

As before, let $[\vec{n}, \vec{u}] = 0$

Also, let $\vec{n}$ be spatial in Frame of A: $\vec{n} \cdot \vec{u} = 0$

Also, let A & B be parallel at point P: $\nabla_{\vec{n}} \vec{u} = 0$ at P.

Finally, A & B have gyroscopes, pointing along $\vec{\sigma}$.

$\vec{\sigma}$ is spatial: $\vec{\sigma} \cdot \vec{u} = 0$

$\vec{\sigma}$ is FW Transported: $\nabla_{\vec{u}} \vec{\sigma} = (\vec{a} \cdot \vec{\sigma}) \vec{u} - (\vec{u} \cdot \vec{\sigma}) \vec{a}$

Difference between $\vec{\sigma}$ at Q and P is $\nabla_{\vec{n}} \vec{\sigma}$

How does $\nabla_{\vec{n}} \vec{\sigma}$ change along world line?

Time rate of change of $\vec{\sigma}$ at Q, as measured by observer at P.

i.e. we want $\nabla_{\vec{u}} \nabla_{\vec{n}} \vec{\sigma}$.

$$\begin{aligned} \nabla_{\vec{u}} \nabla_{\vec{n}} \vec{\sigma} &= \nabla_{\vec{n}} \nabla_{\vec{u}} \vec{\sigma} + \tilde{R}(-, \vec{\sigma}, \vec{u}, \vec{n}) \\ &= \vec{u} \nabla_{\vec{n}} (\vec{a} \cdot \vec{\sigma}) + \tilde{R}(-, \vec{\sigma}, \vec{u}, \vec{n}) \end{aligned}$$

$$\text{or } (\nabla_{\vec{u}} \nabla_{\vec{n}} \vec{\sigma})^\alpha = R^\alpha{}_{\epsilon\gamma\delta} \sigma^\epsilon u^\gamma n^\delta + u^\alpha \nabla_{\vec{n}} (\vec{\sigma} \cdot \vec{\sigma})$$

$$\text{Spatial part of this is } P^\alpha{}_\lambda R^\lambda{}_{\epsilon\gamma\delta} \sigma^\epsilon u^\gamma n^\delta$$

$$\text{where } P^\alpha{}_\lambda = \delta^\alpha{}_\lambda + u^\alpha u_\lambda \quad \text{Projection operator}$$

$$P^\alpha{}_\lambda u^\lambda = P^\alpha{}_\lambda u^\lambda = 0$$

$$\text{Let's call this } P^\alpha{}_\lambda (\nabla_{\vec{u}} \nabla_{\vec{n}} \vec{\sigma})^\lambda = \dot{\sigma}^\alpha = P^\alpha{}_\lambda R^\lambda{}_{\epsilon\gamma\delta} \sigma^\epsilon u^\gamma n^\delta$$

for simplicity, go to rest frame.

$$\text{Note } \dot{\sigma}^\alpha \sigma_\alpha = 0 \\ \dot{\sigma}^\alpha u_\alpha = 0$$

$$\text{so write } \dot{\sigma}^K = R^K{}_{i0j} \sigma^i n^j \quad \text{in rest frame } u^0 = 1$$

$$\text{Let } B_{ij} = \frac{1}{2} \epsilon_{ipq} R^{pq}{}_{j0} \quad \text{"Magnetic part of Riemann tensor"}$$

"Frame-drag Field"

$$\text{Then } \epsilon^K{}_{ps} B^s{}_q = R^K{}_{pq0} = -R^K{}_{p0q}$$

$$\Rightarrow \dot{\sigma}^K = \epsilon^K{}_{si} (B^s{}_j n^j) \sigma^i$$

$$\text{compare with } \frac{d\underline{v}}{dt} = \underline{\Omega} \times \underline{v} \quad \text{Rotation w Frequency } \underline{\Omega}$$

$$\text{so } \dot{\underline{\sigma}} = \underline{\Omega} \times \underline{\sigma} \quad \text{where } \boxed{\Omega^i = B^i{}_j n^j}$$

Gyroscope at Q precesses relative to gyroscope at P.

Bianchi Identity

$$R^\alpha{}_{\epsilon[\delta\gamma;\epsilon]} = 0$$

Why? in LLF $R^\alpha{}_{\epsilon\delta\gamma} = \Gamma^\alpha{}_{\delta\gamma,\epsilon} - \Gamma^\alpha{}_{\epsilon\gamma,\delta}$

and $R^\alpha{}_{\epsilon\delta\gamma;\epsilon} = R^\alpha{}_{\epsilon\delta\gamma,\epsilon}$

$$\begin{aligned} \text{so } R^\alpha{}_{\epsilon[\delta\gamma;\epsilon]} &= \cancel{\Gamma^\alpha{}_{\delta\gamma,\epsilon}} - \cancel{\Gamma^\alpha{}_{\epsilon\gamma,\delta}} \\ &\quad + \cancel{\Gamma^\alpha{}_{\epsilon\delta,\gamma}} - \cancel{\Gamma^\alpha{}_{\delta\epsilon,\gamma}} \\ &\quad + \cancel{\Gamma^\alpha{}_{\epsilon\gamma,\delta}} - \cancel{\Gamma^\alpha{}_{\delta\gamma,\epsilon}} = 0 \end{aligned}$$

Will be important later...