

Significance of LLF: Equivalence principle

Weak equivalence principle:

Motion of Freely falling particles are the same in a grav. Field and in accelerated Frame in Flat spacetime, in small regions.

Einstein equivalence principle:

In small enough regions of spacetime, the non gravitational laws of physics reduce to SR. (ie. cannot detect gravitational field by local experiment)

Strong equivalence principle:

Including gravitational physics.
and particle w/ self gravity, acts same as other particles.

Local experiments $\rightarrow$ LLF

LLF is anfreely-falling frame, $\rightarrow$ instantaneously or local rest Frame.

So in a particle's own LLF, e.g. $U^0 = (1, 0)$

so $\vec{U} \cdot \vec{U} = -1$ in LLF

$\Rightarrow \vec{U} \cdot \vec{U} = -1$ in all coordinate systems.

Commutators of vectors

vector = directional deriv.

Consider two vector fields $\vec{u}, \vec{v}$

$$\vec{u} = \partial_{\vec{u}}$$

$$\vec{v} = \partial_{\vec{v}}$$

Define $[\vec{u}, \vec{v}] = \partial_{\vec{u}} \partial_{\vec{v}} - \partial_{\vec{v}} \partial_{\vec{u}}$

looks tensorish. But this is a vector!

Huh? Proof in coord basis:

$$\begin{aligned} [\vec{u}, \vec{v}] &= \left(u^\alpha \frac{\partial}{\partial x^\alpha} \right) \left(v^\beta \frac{\partial}{\partial x^\beta} \right) - \left(v^\beta \frac{\partial}{\partial x^\beta} \right) \left(u^\alpha \frac{\partial}{\partial x^\alpha} \right) \\ &= u^\alpha v^\beta \frac{\partial^2}{\partial x^\alpha \partial x^\beta} + u^\alpha \frac{\partial v^\beta}{\partial x^\alpha} \frac{\partial}{\partial x^\beta} - v^\beta u^\alpha \frac{\partial^2}{\partial x^\beta \partial x^\alpha} - v^\beta \frac{\partial u^\alpha}{\partial x^\beta} \frac{\partial}{\partial x^\alpha} \\ &= u^\alpha v^\beta_{,\alpha} \frac{\partial}{\partial x^\beta} - v^\alpha u^\beta_{,\alpha} \frac{\partial}{\partial x^\beta} \\ &= \underbrace{(u^\alpha v^\beta_{,\alpha} - v^\alpha u^\beta_{,\alpha})}_{\text{component}} \underbrace{\frac{\partial}{\partial x^\beta}}_{\text{basis vector}} \Rightarrow \text{a vector} \end{aligned}$$

Note $\left[\frac{\partial}{\partial x^\alpha}, \frac{\partial}{\partial x^\beta} \right] = 0$ ("holonomic")

Can be thought of as defn of coord basis: $[\vec{e}_\alpha, \vec{e}_\beta] = 0$

For non-coord basis $[\vec{e}_\alpha, \vec{e}_\beta] \neq 0$
(nonholonomic)

but it is still a vector, so it can be written as $V^\gamma \vec{e}_\gamma$

Define commutation coefficients $C_{\alpha\beta}^\gamma$ by $[\vec{e}_\alpha, \vec{e}_\beta] = C_{\alpha\beta}^\gamma \vec{e}_\gamma$

$C_{\alpha\beta}^\gamma$ has 3 indices, but is not a rank-3 tensor.

($C_{\alpha\beta}^\gamma$ are γ components of a vector resulting from $[\vec{e}_\alpha, \vec{e}_\beta]$)

In a general basis,

$$[\vec{u}, \vec{v}] = [u^\alpha \vec{e}_\alpha, v^\beta \vec{e}_\beta]$$

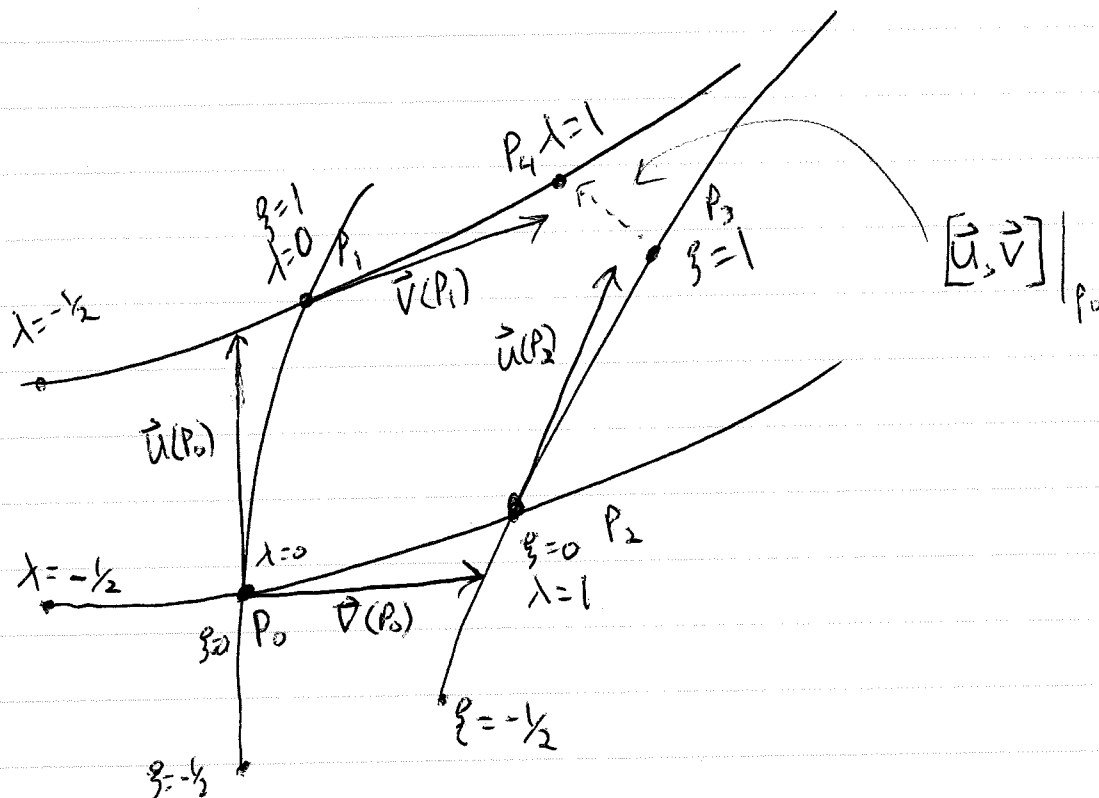
$$= u^\alpha \vec{e}_\alpha v^\beta \vec{e}_\beta - v^\beta \vec{e}_\beta u^\alpha \vec{e}_\alpha$$

$$= u^\alpha (\partial_\alpha v^\beta) \vec{e}_\beta + u^\alpha v^\beta \vec{e}_\alpha \vec{e}_\beta - v^\beta (\partial_\beta u^\alpha) \vec{e}_\alpha - v^\beta u^\alpha \vec{e}_\beta \vec{e}_\alpha$$

$$= u^\alpha v^\beta \underbrace{[\vec{e}_\alpha, \vec{e}_\beta]}_{C_{\alpha\beta}^\gamma \vec{e}_\gamma} + u^\alpha (\partial_\alpha v^\gamma) \vec{e}_\gamma - v^\beta (\partial_\beta u^\gamma) \vec{e}_\gamma$$

$$= (u^\alpha v^\beta + u^\alpha \partial_\alpha v^\gamma - v^\beta \partial_\beta u^\gamma) \vec{e}_\gamma$$

Picture of commutator



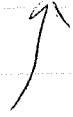
$$\vec{U} = \frac{d}{d\beta} \quad \vec{V} = \frac{d}{d\lambda}$$

Consider scalar function f , in coord basis

$$\begin{aligned} f(P_4) - f(P_3) &= f(P_4) - f(P_1) + f(P_1) - f(P_0) \\ &\quad + f(P_0) - f(P_2) + f(P_2) - f(P_3) \end{aligned}$$

$$\begin{aligned} &= \left(f_{,\alpha} V^\alpha + \frac{1}{2} f_{,\alpha\beta} V^\alpha V^\beta \right)_{P_1} + \left(f_{,\alpha} U^\alpha + \frac{1}{2} f_{,\alpha\beta} U^\alpha U^\beta \right)_{P_0} \\ &\quad - \left(f_{,\alpha} V^\alpha + \frac{1}{2} f_{,\alpha\beta} V^\alpha V^\beta \right)_{P_0} - \left(f_{,\alpha} U^\alpha + \frac{1}{2} f_{,\alpha\beta} U^\alpha U^\beta \right)_{P_2} \end{aligned}$$

$$= u^\alpha \partial_\alpha \left[b_{\epsilon} v^\epsilon + \frac{1}{2} b_{\epsilon\gamma} v^\epsilon v^\gamma \right] - v^\alpha \partial_\alpha \left[b_{\epsilon} u^\epsilon + \frac{1}{2} b_{\epsilon\gamma} u^\epsilon u^\gamma \right]$$



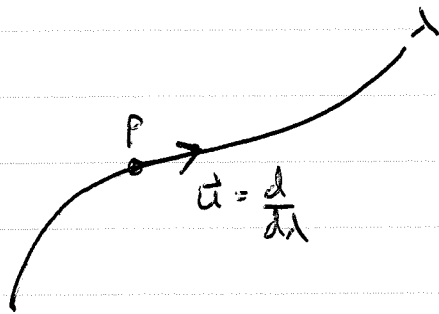
small when $u, v \rightarrow 0$

$$= u^\alpha \partial_\alpha [v^\epsilon \partial_\epsilon f] - v^\alpha \partial_\alpha [u^\epsilon \partial_\epsilon f]$$

$$= [u^\alpha \partial_\alpha v^\epsilon - v^\alpha \partial_\alpha u^\epsilon] \partial_\epsilon f$$

$$= [\vec{u}, \vec{v}]^\epsilon \partial_\epsilon f$$

Covariant Derivative



We know how to differentiate scalar fields along a curve

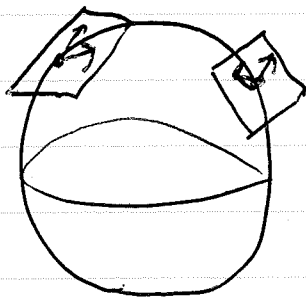
$$\frac{d}{d\lambda} \phi = \partial_{\vec{u}} \phi = u^\alpha \partial_\alpha \phi$$

$$= u^\alpha \frac{\partial \phi}{\partial x^\alpha} \quad \text{in coord basis.}$$

$$= \lim_{\epsilon \rightarrow 0} \frac{\phi(\lambda + \epsilon) - \phi(\lambda)}{\epsilon}$$

← compare ϕ at nearby pts along curve

But how to differentiate vectors along a curve?



Vectors at different points live in different vector spaces.

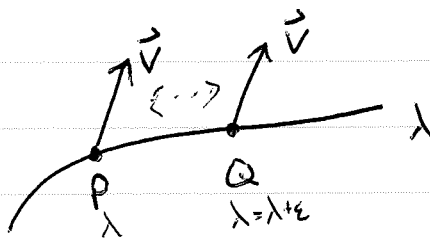
⇒ cannot naturally compare vectors at different points.

So we need a rule for comparing vectors at nearby points.
↓
+ tensors

Rule we will use is parallel transport:

Two derivations

① Physical.



Go to LLF at point P. Parallel transport vector at Q to vector at P by keeping components of $\vec{v}$ unchanged in LLF.

$$\text{Then } \frac{d}{d\lambda} \vec{V} \Big|_{\lambda_0} = \lim_{\epsilon \rightarrow 0} \frac{\vec{V}(\lambda_0 + \epsilon) \text{ parallel transported to } \lambda_0 - \vec{V}(\lambda_0)}{\epsilon}$$

Write this as $\nabla_{\vec{u}} \vec{V}$ where $\vec{u} = \frac{d}{d\lambda}$
 ↑ covariant derivative along $\vec{u}$

Similarly for tensors

$$\text{let } \tilde{T} = T^\alpha_{\beta} \vec{e}_\alpha \otimes \tilde{\omega}^\beta \quad \text{in LLF}$$

$$\nabla_{\vec{u}} \tilde{T} = \lim_{\epsilon \rightarrow 0} \frac{T^\alpha_{\beta} \vec{e}_\alpha \otimes \tilde{\omega}^\beta \Big|_{\epsilon \text{ to } 0}^{\text{transport from } \epsilon \text{ to } 0} - T^\alpha_{\beta} \vec{e}_\alpha \otimes \tilde{\omega}^\beta \Big|_{\epsilon \rightarrow 0}}{\epsilon}$$

but basis vectors constant in LLF

$$\text{So } \nabla_{\vec{u}} \tilde{T} = u^\gamma T^\alpha_{\beta, \gamma} \vec{e}_\alpha \otimes \tilde{\omega}^\beta \quad \text{In LLF}$$

② More abstract defn of cov. derivative.

Define $\nabla \tilde{T}$ as a tensor

covariant derivative of $\tilde{T}$

such that $\nabla \tilde{T}(\dots, \vec{u}) = \nabla_{\vec{u}} \tilde{T}$

↑ slots of $\tilde{T}$

(If $\tilde{T}$ has components T^α_β , $\nabla \tilde{T}$ has components $\nabla_\gamma T^\alpha_\beta$
 $\nabla_{\vec{u}} \tilde{T}$ has components $u^\gamma \nabla_\gamma T^\alpha_\beta$
 $\nabla_\gamma T^\alpha_\beta$ also written $T^\alpha_{\beta;\gamma}$)

and such that $\nabla \equiv$ covariant derivative operator has following properties

① Linear: $\nabla(\alpha \tilde{T} + \beta \tilde{S}) = \alpha \nabla \tilde{T} + \beta \nabla \tilde{S}$
 α, β constants

② Leibnitz: $\nabla_{\vec{u}}(f \tilde{A} \otimes \tilde{B}) = (\nabla_{\vec{u}} f) \tilde{A} \otimes \tilde{B}$
 $+ f \nabla_{\vec{u}} \tilde{A} \otimes \tilde{B}$
 $+ f \tilde{A} \otimes \nabla_{\vec{u}} \tilde{B}$

③ Commutes w/ contraction: if $\tilde{S}$ is rank n ,
 $\tilde{S}(\vec{e}_\alpha, \vec{\omega}^\alpha, \dots)$ is rank $n-2$

$\nabla[\tilde{S}(\vec{e}_\alpha, \vec{\omega}^\alpha, \dots)] = \underbrace{\nabla \tilde{S}(\vec{e}_\alpha, \vec{\omega}^\alpha, \dots)}_{\text{rank } n+1}$

Or

$$S^\alpha_{\alpha\beta;\gamma}$$

doesn't depend on if slots
contracted before/after deriv

① For scalars,

$$\begin{aligned}\partial_{\vec{u}} \phi &= \nabla_{\vec{u}} \phi = u^\alpha \nabla_\alpha \phi = u^\alpha \phi_{,\alpha} \\ &= u^\alpha \partial_\alpha \phi = u^\alpha \phi_{,\alpha}\end{aligned}$$

$$\textcircled{e} \quad \nabla_{\vec{u}} \vec{v} - \nabla_{\vec{v}} \vec{u} = [\vec{u}, \vec{v}]$$

$$\left(\text{equivalent to } (\nabla \nabla f)(\vec{u}, \vec{v}) = (\nabla \nabla f)(\vec{v}, \vec{u}) \right)$$

torsion free

②

$$\nabla_{\vec{u}} \tilde{g} = 0$$

compatibility
of covariant derivative
w/ metric